

Relation between \underline{E} and \underline{H} in a uniform plane wave $\nabla \cdot \underline{E} = \frac{1}{\epsilon} \nabla \cdot \underline{D} = 0$

$$\Rightarrow \frac{\partial E_x}{\partial x} + \frac{\partial E_y}{\partial y} + \frac{\partial E_z}{\partial z} = 0 \quad \text{no variation of } E_z \text{ in the } z \text{ direction (direction of propagation)}$$

$$\Rightarrow \frac{\partial E_z}{\partial z} = 0 \Rightarrow E_z \text{ be either 0 or constant in space .}$$

$$\Rightarrow \frac{\partial^2 E_z}{\partial t^2} = 0 \Rightarrow E_z \text{ can at most be a linearly increasing function of time. (Inna, p643)}$$

$$\Rightarrow E_z = 0 \quad \therefore \text{ a uniform plane wave propagating in the } z \text{ direction has no } z \text{ component of}$$

Assuming only E_x exists

$$\nabla \times \underline{E} = -\mu_0 \frac{\partial \underline{H}}{\partial t}$$

$$\nabla \times \underline{E} = \begin{vmatrix} \underline{a}_x & \underline{a}_y & \underline{a}_z \\ 0 & 0 & \frac{\partial}{\partial z} \\ E_x & 0 & 0 \end{vmatrix} = \frac{\partial E_x}{\partial z} \underline{a}_y = -\mu_0 \frac{\partial H_y}{\partial t} \underline{a}_y$$

$$\frac{\partial H_y}{\partial t} = -\frac{1}{\mu_0} \frac{\partial E_x}{\partial z}$$

$$H_y = -\frac{1}{\mu_0} \int \frac{\partial E_x}{\partial z} dt + C$$

$$\text{if } E_x = E_x(z - v_p t), \quad \zeta = z - v_p t$$

$$\text{then } \begin{cases} \frac{\partial E_x(z - v_p t)}{\partial z} = \frac{\partial E_x(\zeta)}{\partial \zeta} \frac{\partial \zeta}{\partial z} = \frac{\partial E_x(\zeta)}{\partial \zeta} \\ \frac{d\zeta}{dt} = -v_p, \quad dt = -\frac{d\zeta}{v_p} \end{cases}$$

$$\begin{aligned} H_y &= -\frac{1}{\mu_0} \int \frac{\partial E_x}{\partial z} dt + C \\ &= -\frac{1}{\mu_0} \int \frac{\partial E_x(\zeta)}{\partial \zeta} \left(-\frac{d\zeta}{v_p} \right) + C \\ &= \frac{1}{\mu_0 v_p} E_x(\zeta) + C \end{aligned}$$

$$H_y = \frac{1}{\mu_0 v_p} E_x \quad \text{in free space, } v_p = 1/\sqrt{\mu_0 \epsilon_0}$$

C indicates that a field could have component that is independent of t which would not be a part of wave motion ($z - v_p t$).

$$\frac{E_x}{H_y} = \sqrt{\frac{\mu_0}{\epsilon_0}} = \eta_0 = 120\pi = 377 \quad [\Omega] \quad \text{intrinsic impedance of the free space}$$

$$\text{intrinsic impedance of an medium } \eta \equiv \sqrt{\mu/\epsilon} \quad [\Omega]$$

Similarly,

$$\frac{E_y}{H_x} = -\sqrt{\frac{\mu_0}{\epsilon_0}} = -\eta_0$$

$$\underline{E} \cdot \underline{H} = E_x H_x + E_y H_y = \eta_0 H_y H_x - \eta_0 H_x H_y = 0$$

$$\therefore \underline{E} \perp \underline{H}$$

$$\underline{E} \times \underline{H} = \begin{vmatrix} \underline{a}_x & \underline{a}_y & \underline{a}_z \\ E_x & E_y & 0 \\ H_x & H_y & 0 \end{vmatrix} = (E_x H_y - E_y H_x) \underline{a}_z$$

Electromagnetic energy flow and the Poynting vector

$$\underline{F} = q(\underline{E} + \underline{v} \times \underline{B})$$

$$dW = q(\underline{E} + \underline{v} \times \underline{B}) \cdot d\underline{l}; \quad d\underline{l} = \underline{v} dt, \quad dW = P dt$$

$$P = q(\underline{E} + \underline{v} \times \underline{B}) \cdot d\underline{v}$$

$$= q \underline{v} \cdot \underline{E}$$

$$dP = \rho_v d\underline{v} \cdot \underline{E} \quad q \rightarrow \rho_v d\underline{v}$$

$$= (\rho_v \underline{v}) d\underline{v} \cdot \underline{E}$$

$$= \underline{J} d\underline{v} \cdot \underline{E} \quad \underline{J} = \rho_v \underline{v}$$

$$\text{power density } p = \frac{dP}{d\underline{v}} = \underline{J} \cdot \underline{E}$$

$$\nabla \times \underline{H} = \underline{J} + \frac{\partial \underline{D}}{\partial t}, \quad \underline{J} = \nabla \times \underline{H} - \frac{\partial \underline{D}}{\partial t}$$

$$\underline{J} \cdot \underline{E} = \underline{E} \cdot (\nabla \times \underline{H}) - \underline{E} \cdot \frac{\partial \underline{D}}{\partial t}$$

$$\nabla \cdot (\underline{E} \times \underline{H}) = \underline{H} \cdot (\nabla \times \underline{E}) - \underline{E} \cdot (\nabla \times \underline{H})$$

$$\underline{J} \cdot \underline{E} = \underline{H} \cdot (\nabla \times \underline{E}) - \nabla \cdot (\underline{E} \times \underline{H}) - \underline{E} \cdot \frac{\partial \underline{D}}{\partial t}$$

$$\underline{J} \cdot \underline{E} = -\underline{H} \cdot \frac{\partial \underline{B}}{\partial t} - \underline{E} \cdot \frac{\partial \underline{D}}{\partial t} - \nabla \cdot (\underline{E} \times \underline{H}), \quad \nabla \times \underline{E} = -\frac{\partial \underline{B}}{\partial t}$$

$$\text{Poynting vector } \underline{S} = \underline{E} \times \underline{H} \quad [\text{W}/\text{m}^2]$$

for the time-varying fields in a linear, homogeneous, and isotropic medium, $\underline{B} = \mu \underline{H}, \underline{D} = \epsilon \underline{E}$

$$\underline{H} \cdot \frac{\partial \underline{B}}{\partial t} = \frac{1}{2} \frac{\partial}{\partial t} (\underline{B} \cdot \underline{H}) = \frac{\partial}{\partial t} \left(\frac{1}{2} \mu |\underline{H}|^2 \right)$$

$$\underline{E} \cdot \frac{\partial \underline{D}}{\partial t} = \frac{1}{2} \frac{\partial}{\partial t} (\underline{D} \cdot \underline{E}) = \frac{\partial}{\partial t} \left(\frac{1}{2} \epsilon |\underline{E}|^2 \right)$$

$$\underline{J} \cdot \underline{E} = -\frac{\partial}{\partial t} \left(\frac{1}{2} \mu |\underline{H}|^2 \right) - \frac{\partial}{\partial t} \left(\frac{1}{2} \epsilon |\underline{E}|^2 \right) - \nabla \cdot (\underline{E} \times \underline{H}) \quad \text{differential form of Poynting's theorem}$$

theorem

$$\int_v \underline{J} \cdot \underline{E} dv = -\frac{\partial}{\partial t} \int_v \left(\frac{1}{2} \mu |\underline{H}|^2 + \frac{1}{2} \epsilon |\underline{E}|^2 \right) dv - \int_v \nabla \cdot (\underline{E} \times \underline{H}) dv$$

Poynting theorem $-\int_v (\underline{E} \times \underline{H}) \cdot d\underline{s} = \int_v \underline{J} \cdot \underline{E} dv + \frac{\partial}{\partial t} \int_v \left(\frac{1}{2} \mu |\underline{H}|^2 + \frac{1}{2} \epsilon |\underline{E}|^2 \right) dv$ conservation of energy

Electromagnetic power flow into a closed surface at any instant equals the sum of the time rates of increase of the stored electric and magnetic energies plus the ohmic power dissipated (or electric power generated, if the surface enclosed a source) within the enclosed volume.

$\int_v (\underline{E} \times \underline{H}) \cdot d\underline{s}$ the flow of energy inward or outward through the surface enclosing the volume v

$\int_v \underline{E} \cdot \underline{J} dv$ Joule's law : the instantaneous power dissipated in the volume v (ohmic power loss)

$\frac{\partial}{\partial t} \int_v \left(\frac{1}{2} \mu |\underline{H}|^2 + \frac{1}{2} \epsilon |\underline{E}|^2 \right) dv$ the rate of decrease of the electromagnetic energy stored in volume v

electric energy density $w_e = \frac{1}{2} \epsilon |\underline{E}|^2 \quad [J / m^3]$

magnetic energy density $w_m = \frac{1}{2} \mu |\underline{H}|^2 \quad [J / m^3]$

$$-\int_v \underline{S} \cdot d\underline{s} = \int_v \underline{J} \cdot \underline{E} dv + \frac{\partial}{\partial t} \int_v (w_e + w_m) dv$$

Time-harmonic uniform plane waves in a lossless medium

$$\underline{E}(x, y, z, t) = \underline{E}(\underline{r}, t) = \text{Re} \{ \hat{\underline{E}}(\underline{r}) e^{j\omega t} \} = \frac{1}{2} \left(\hat{\underline{E}}(\underline{r}) e^{j\omega t} + \hat{\underline{E}}^*(\underline{r}) e^{-j\omega t} \right)$$

$$\underline{H}(x, y, z, t) = \underline{H}(\underline{r}, t) = \text{Re} \{ \hat{\underline{H}}(\underline{r}) e^{j\omega t} \} = \frac{1}{2} \left(\hat{\underline{H}}(\underline{r}) e^{j\omega t} + \hat{\underline{H}}^*(\underline{r}) e^{-j\omega t} \right)$$

$$\underline{J}(x, y, z, t) = \underline{J}(\underline{r}, t) = \text{Re} \{ \hat{\underline{J}}(\underline{r}) e^{j\omega t} \}$$

$$\rho(x, y, z, t) = \rho(\underline{r}, t) = \text{Re} \{ \hat{\rho}(\underline{r}) e^{j\omega t} \}$$

$$\frac{\partial}{\partial t} \rightarrow j\omega$$

$$\nabla \times \underline{\hat{E}} = -j\omega\mu\underline{\hat{H}} \qquad \nabla \times \underline{\hat{H}} = j\omega\varepsilon\underline{\hat{E}}$$

$$\nabla \cdot \underline{\hat{D}} = \rho \qquad \nabla \cdot \underline{\hat{B}} = 0$$

$$\nabla \times \nabla \times \underline{\hat{E}} = -j\omega\mu\nabla \times \underline{\hat{H}}$$

$$\nabla(\nabla \cdot \underline{\hat{E}}) - \nabla^2 \underline{\hat{E}} = -j\omega\mu(-j\omega\varepsilon\underline{\hat{E}})$$

$$\nabla^2 \underline{\hat{E}} + \omega^2 \mu\varepsilon \underline{\hat{E}} = 0, \qquad \nabla^2 \underline{\hat{H}} + k^2 \underline{\hat{H}} = 0$$

$$k = \omega\sqrt{\mu\varepsilon} \quad \text{phase constant; wave number; propagation constant}$$

$$\nabla \times \nabla \times \underline{\hat{H}} = j\omega\varepsilon\nabla \times \underline{\hat{E}}$$

$$\nabla(\nabla \cdot \underline{\hat{H}}) - \nabla^2 \underline{\hat{H}} = j\omega\varepsilon(-j\omega\mu\underline{\hat{H}})$$

$$\nabla^2 \underline{\hat{H}} + \omega^2 \mu\varepsilon \underline{\hat{H}} = 0, \qquad \nabla^2 \underline{\hat{H}} + k^2 \underline{\hat{H}} = 0$$

$$\underline{\hat{E}} = \hat{E}_x \underline{a}_x + \hat{E}_y \underline{a}_y + \hat{E}_z \underline{a}_z$$

$$\left[\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right] \hat{E}_x + k^2 \hat{E}_x = 0$$

$$\left[\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right] \hat{E}_y + k^2 \hat{E}_y = 0$$

$$\left[\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right] \hat{E}_z + k^2 \hat{E}_z = 0$$

Considering only \hat{E}_x

$$\text{Let } \hat{E}_x(x, y, z) = \hat{f}(x)\hat{g}(y)\hat{h}(z)$$

$$\left[\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right] \hat{E}_x + k^2 \hat{E}_x = 0$$

$$\hat{f}''\hat{g}\hat{h} + \hat{f}\hat{g}''\hat{h} + \hat{f}\hat{g}\hat{h}'' + k^2 \hat{f}\hat{g}\hat{h} = 0$$

$$\frac{\hat{f}''}{\hat{f}} + \frac{\hat{g}''}{\hat{g}} + \frac{\hat{h}''}{\hat{h}} + k^2 = 0$$

$$\begin{cases} \frac{\hat{f}''}{\hat{f}} = -k_x^2 & \hat{f}'' + k_x^2 \hat{f} = 0 \\ \frac{\hat{g}''}{\hat{g}} = -k_y^2 & \hat{g}'' + k_y^2 \hat{g} = 0 \\ \frac{\hat{h}''}{\hat{h}} = -k_z^2 & \hat{h}'' + k_z^2 \hat{h} = 0 \end{cases}$$

$$\begin{aligned}\hat{f} &= Ae^{-jk_x x} + Be^{jk_x x} \\ \Rightarrow \hat{g} &= Ce^{-jk_y y} + De^{jk_y y} \\ \hat{h} &= Ee^{-jk_z z} + Fe^{jk_z z}\end{aligned}$$

Assuming only $e^{-jk_x x}$ term exists

$$\begin{aligned}\hat{E}_x(x, y, z) &= \hat{f}(x)\hat{g}(y)\hat{h}(z) = ACEe^{-j(k_x x + k_y y + k_z z)} \\ \Rightarrow \hat{E}_x(x, y, z) &= \hat{E}_{0x}e^{-j(k_x x + k_y y + k_z z)}\end{aligned}$$

$$\Rightarrow \begin{cases} \hat{E}_x(x, y, z) = \hat{E}_{0x}e^{-j(k_x x + k_y y + k_z z)} \\ \hat{E}_y(x, y, z) = \hat{E}_{0y}e^{-j(k_x x + k_y y + k_z z)} \\ \hat{E}_z(x, y, z) = \hat{E}_{0z}e^{-j(k_x x + k_y y + k_z z)} \end{cases} \quad \hat{E}_{0x}, \hat{E}_{0y}, \hat{E}_{0z} \text{ are constant}$$

$$\underline{k} = k_x \underline{a}_x + k_y \underline{a}_y + k_z \underline{a}_z = k \underline{a}_n, \quad k = \sqrt{k_x^2 + k_y^2 + k_z^2}$$

$$\underline{r} = x \underline{a}_x + y \underline{a}_y + z \underline{a}_z$$

$$\underline{k} \cdot \underline{r} = k_x x + k_y y + k_z z$$

$$\begin{aligned}\hat{E}(x, y, z) &= \hat{E}_x \underline{a}_x + \hat{E}_y \underline{a}_y + \hat{E}_z \underline{a}_z \\ &= \left[\hat{E}_{0x} \underline{a}_x + \hat{E}_{0y} \underline{a}_y + \hat{E}_{0z} \underline{a}_z \right] e^{-j(k_x x + k_y y + k_z z)} \quad \underline{\hat{E}}_0 \text{ is a constant vector} \\ &= \underline{\hat{E}}_0 e^{-j \underline{k} \cdot \underline{r}}\end{aligned}$$

$$\begin{aligned}\underline{E}(\underline{r}, t) &= \text{Re} \left\{ \hat{E}(x, y, z) e^{j\omega t} \right\} \\ &= \text{Re} \left\{ \underline{\hat{E}}_0 e^{j(\omega t - \underline{k} \cdot \underline{r})} \right\}\end{aligned}$$

Generally speaking, for $\nabla^2 \underline{\hat{E}} + k^2 \underline{\hat{E}} = 0$

$$\underline{\hat{E}}(\underline{r}) = \hat{E}(x, y, z) = \underline{\hat{E}}_0^+ e^{-j \underline{k} \cdot \underline{r}} + \underline{\hat{E}}_0^- e^{j \underline{k} \cdot \underline{r}}$$

$\underline{\hat{E}}(x, y, z) = \hat{E}_x(z) \underline{a}_x$ of a uniform plane wave is uniform over the xy plane

\Rightarrow independent of the x and y coordinates

$$\Rightarrow \frac{\partial \hat{E}_x}{\partial x} = \frac{\partial \hat{E}_x}{\partial y} = 0$$

$$\frac{\partial^2 \hat{E}_x}{\partial z^2} + k^2 \hat{E}_x = 0$$

$$\hat{E}_x(z) = \underbrace{\hat{E}_{x0}^+ e^{-jkz}}_{\hat{E}_x^+(z)} + \underbrace{\hat{E}_{x0}^- e^{jkz}}_{\hat{E}_x^-(z)}$$

$$E_x(z, t) = \text{Re} \left\{ \left(\hat{E}_{x0}^+ e^{-jkz} + \hat{E}_{x0}^- e^{jkz} \right) e^{j\omega t} \right\}$$

Relation between $\underline{\hat{E}}$ and $\underline{\hat{H}}$ in a uniform plane wave

$$\begin{aligned} \nabla \cdot \underline{\hat{E}} = 0 &= \nabla \cdot \underline{\hat{E}}_0 e^{-jk \cdot r} \\ &= e^{-jk \cdot r} \nabla \cdot \underline{\hat{E}}_0 + \underline{\hat{E}}_0 \cdot \nabla e^{-jk \cdot r} \\ &= \underline{\hat{E}}_0 \cdot \nabla e^{-jk \cdot r} \\ &= \underline{\hat{E}}_0 \cdot [-j\mathbf{k} e^{-jk \cdot r}] \\ &= -j(\mathbf{k} \cdot \underline{\hat{E}}_0 e^{-jk \cdot r}) \\ &= -j(\mathbf{k} \cdot \underline{\hat{E}}) \end{aligned}$$

$$\begin{aligned} \nabla e^{-jk \cdot r} &= \left(\frac{\partial}{\partial x} \underline{a}_x + \frac{\partial}{\partial y} \underline{a}_y + \frac{\partial}{\partial z} \underline{a}_z \right) e^{-j(k_x x + k_y y + k_z z)} \\ &= -j(k_x \underline{a}_x + k_y \underline{a}_y + k_z \underline{a}_z) e^{-j(k_x x + k_y y + k_z z)} \\ &= -j\mathbf{k} e^{-j(k \cdot r)} \\ &= -j\mathbf{k} e^{-j(k \cdot r)} \quad \nabla \rightarrow \pm j\mathbf{k} \end{aligned}$$

$$\mathbf{k} \cdot \underline{\hat{E}} = 0 \quad \Rightarrow \quad \mathbf{k} \perp \underline{\hat{E}} \quad \Rightarrow \quad \mathbf{k} \perp \underline{E}$$

$$\nabla \times \varphi \underline{A} = \varphi \nabla \times \underline{A} - \underline{A} \times \nabla \varphi$$

$$\nabla \cdot \varphi \underline{A} = \varphi \nabla \cdot \underline{A} + \underline{A} \cdot \nabla \varphi$$

uniform plane wave $\nabla \cdot \underline{\hat{E}}_0 = 0$

$$\begin{aligned} -j\omega\mu_0 \underline{\hat{H}} &= \nabla \times \underline{\hat{E}} = \nabla \times \underline{\hat{E}}_0 e^{-jk \cdot r} \\ -j\omega\mu_0 \underline{\hat{H}} &= -\underline{\hat{E}}_0 \times \nabla e^{-jk \cdot r} \end{aligned}$$

$$\begin{aligned} j\omega\varepsilon_0 \underline{\hat{E}} &= \nabla \times \underline{\hat{H}} = \nabla \times \underline{\hat{H}}_0 e^{-jk \cdot r} \\ j\omega\varepsilon_0 \underline{\hat{E}} &= -\underline{\hat{H}}_0 \times \nabla e^{-jk \cdot r} \end{aligned}$$

$$\begin{aligned} \underline{\hat{H}} &= -\frac{1}{j\omega\mu_0} \underline{\hat{E}}_0 \times -j\mathbf{k} e^{-jk \cdot r} \\ &= \frac{\mathbf{k} \times \underline{\hat{E}}_0 e^{-jk \cdot r}}{\omega\mu_0} = \frac{\mathbf{k} \times \underline{\hat{E}}}{\omega\mu_0} \\ &= \frac{k}{\omega\mu_0} \underline{a}_n \times \underline{\hat{E}}_0 e^{-jk \cdot r} \\ &= \frac{1}{\eta_0} \underline{a}_n \times \underline{\hat{E}}_0 e^{-jk \cdot r} = \frac{1}{\eta_0} \underline{a}_n \times \underline{\hat{E}} \\ &= \underline{\hat{H}}_0 e^{-jk \cdot r} \\ \frac{k}{\omega\mu_0} &= \frac{\omega\sqrt{\mu_0\varepsilon_0}}{\omega\mu_0} = \sqrt{\frac{\varepsilon_0}{\mu_0}} = \frac{1}{\eta_0} \end{aligned}$$

$$\begin{aligned} \underline{\hat{E}} &= -\frac{1}{j\omega\varepsilon_0} \underline{\hat{H}}_0 \times -j\mathbf{k} e^{-jk \cdot r} \\ &= \frac{\underline{\hat{H}}_0 e^{-jk \cdot r} \times \mathbf{k}}{\omega\varepsilon_0} = \frac{\underline{\hat{H}} \times \mathbf{k}}{\omega\varepsilon_0} \\ &= \underline{\hat{H}}_0 e^{-jk \cdot r} \times \frac{k}{\omega\varepsilon_0} \underline{a}_n \\ &= \underline{\hat{H}}_0 e^{-jk \cdot r} \times \eta_0 \underline{a}_n = \underline{\hat{H}} \times \eta_0 \underline{a}_n \\ &= \underline{\hat{E}}_0 e^{-jk \cdot r} \\ \frac{k}{\omega\varepsilon_0} &= \frac{\omega\sqrt{\mu_0\varepsilon_0}}{\omega\varepsilon_0} = \sqrt{\frac{\mu_0}{\varepsilon_0}} = \eta_0 \end{aligned}$$

$$\underline{\hat{H}} = \frac{1}{\eta_0} \underline{a}_n \times \underline{\hat{E}}$$

$$\underline{\hat{E}} = \underline{\hat{H}} \times \eta_0 \underline{a}_n$$

$$\eta_0 = \sqrt{\frac{\mu_0}{\epsilon_0}} \quad [\Omega]$$

$$[\mu_0] = [ML/Q^2] \quad [\epsilon_0] = [Q^2T^2/ML^3]$$

$$\Rightarrow \left[\sqrt{\frac{\mu_0}{\epsilon_0}} \right] = \left[\frac{ML/Q^2}{Q^2T^2/ML^3} \right]^{1/2} = \left[\frac{M^2L^4}{Q^4T^2} \right]^{1/2} = \left[\frac{ML^2}{Q^2T} \right]$$

$$[\Omega] = \left[\frac{ML^2}{Q^2T} \right]$$

Poynting Vector

Complex Poynting vector

$$\underline{S} = \underline{E} \times \underline{H} = \frac{1}{2} \left[\hat{\underline{E}}(\underline{r})e^{j\omega t} + \hat{\underline{E}}^*(\underline{r})e^{-j\omega t} \right] \times \frac{1}{2} \left[\hat{\underline{H}}(\underline{r})e^{j\omega t} + \hat{\underline{H}}^*(\underline{r})e^{-j\omega t} \right]$$

$$= \frac{1}{4} \left[\hat{\underline{E}}(\underline{r}) \times \hat{\underline{H}}^*(\underline{r}) + \hat{\underline{E}}^*(\underline{r}) \times \hat{\underline{H}}(\underline{r}) \right] + \frac{1}{4} \left[\hat{\underline{E}}(\underline{r}) \times \hat{\underline{H}}(\underline{r})e^{j2\omega t} + \hat{\underline{E}}^*(\underline{r}) \times \hat{\underline{H}}^*(\underline{r})e^{-j2\omega t} \right]$$

since $\left[\hat{\underline{E}}(\underline{r}) \times \hat{\underline{H}}^*(\underline{r}) \right]^* = \hat{\underline{E}}^*(\underline{r}) \times \hat{\underline{H}}(\underline{r}), \quad \left[\hat{\underline{E}}^*(\underline{r}) \times \hat{\underline{H}}(\underline{r})e^{-j2\omega t} \right]^* = \hat{\underline{E}}(\underline{r}) \times \hat{\underline{H}}(\underline{r})e^{j2\omega t}$

$$\underline{S} = \frac{1}{2} \text{Re} \left\{ \hat{\underline{E}}(\underline{r}) \times \hat{\underline{H}}^*(\underline{r}) \right\} + \frac{1}{2} \text{Re} \left\{ \hat{\underline{E}}(\underline{r}) \times \hat{\underline{H}}(\underline{r})e^{j2\omega t} \right\}$$

time-average Poynting vector $\langle \underline{S}_{av} \rangle = \frac{1}{T} \int_0^T \underline{S} dt = \frac{1}{2} \text{Re} \left\{ \hat{\underline{E}}(\underline{r}) \times \hat{\underline{H}}^*(\underline{r}) \right\}$

complex Poynting vector $\hat{\underline{S}} = \hat{\underline{E}} \times \hat{\underline{H}}^*$

$$\hat{\underline{S}} = \hat{\underline{E}} \times \hat{\underline{H}}^* = \hat{\underline{E}}_0 e^{-jk \cdot \underline{r}} \times \left(\frac{1}{\eta_0} \underline{a}_n \times \hat{\underline{E}}_0 e^{-jk \cdot \underline{r}} \right)^*$$

$$= \hat{\underline{E}}_0 \times \frac{1}{\eta_0} (\underline{a}_n \times \hat{\underline{E}}_0)^* = \frac{1}{\eta_0} \left(\hat{\underline{E}}_0 \cdot \hat{\underline{E}}_0^* \right) \underline{a}_n - \frac{1}{\eta_0} \left(\underbrace{\hat{\underline{E}}_0 \cdot \underline{a}_n}_{=0} \right) \hat{\underline{E}}_0^*$$

$$= \frac{1}{\eta_0} \left| \hat{\underline{E}}_0 \right|^2 \underline{a}_n$$