

Plane of constant phase

Parallel polarization

Oblique incidence at a dielectric boundary

Incident wave	Reflected wave	Refracted wave
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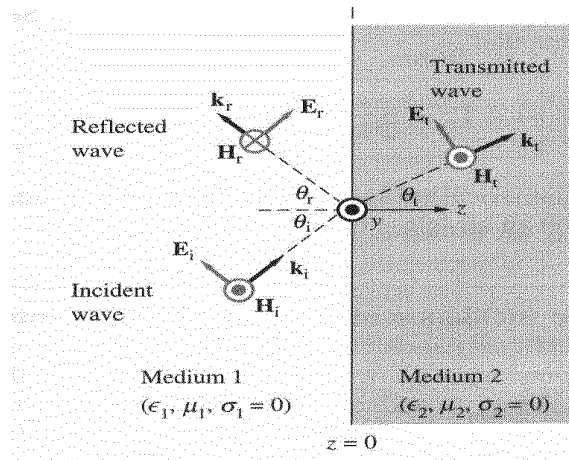
$$\begin{aligned} \underline{\hat{E}}_i &= \hat{E}_{i0} e^{-jk_i \cdot r} (\underline{a}_y \times \underline{a}_{ni}) & \underline{\hat{E}}_r &= \hat{E}_{r0} e^{-jk_r \cdot r} (\underline{a}_y \times \underline{a}_{nr}) & \underline{\hat{E}}_t &= \hat{E}_{t0} e^{-jk_t \cdot r} (\underline{a}_y \times \underline{a}_{nt}) \\ \underline{\hat{H}}_i &= \frac{\hat{E}_{i0} e^{-jk_i \cdot r}}{\eta_1} \underline{a}_y & \underline{\hat{H}}_r &= \frac{\hat{E}_{r0} e^{-jk_r \cdot r}}{\eta_1} \underline{a}_y & \underline{\hat{H}}_t &= \frac{\hat{E}_{t0} e^{-jk_t \cdot r}}{\eta_2} \underline{a}_y \end{aligned}$$

$$\underline{a}_{ni} = \sin \theta_i \underline{a}_x + \cos \theta_i \underline{a}_z \qquad \underline{a}_{nr} = \sin \theta_r \underline{a}_x - \cos \theta_r \underline{a}_z$$

$$\underline{k}_i = k_i \sin \theta_i \underline{a}_x + k_i \cos \theta_i \underline{a}_z \qquad \underline{k}_r = k_r \sin \theta_r \underline{a}_x - k_r \cos \theta_r \underline{a}_z \qquad \underline{k}_t = k_t \sin \theta_t \underline{a}_x + k_t \cos \theta_t \underline{a}_z$$

$$\begin{cases} \underline{\hat{E}}_i = \hat{E}_{i0} (\cos \theta_i \underline{a}_x - \sin \theta_i \underline{a}_z) e^{-j(k_i \sin \theta_i x + k_i \cos \theta_i z)} \\ \underline{\hat{E}}_r = \hat{E}_{r0} (\cos \theta_r \underline{a}_x + \sin \theta_r \underline{a}_z) e^{-j(k_r \sin \theta_r x - k_r \cos \theta_r z)} \\ \underline{\hat{E}}_t = \hat{E}_{t0} (\cos \theta_t \underline{a}_x - \sin \theta_t \underline{a}_z) e^{-j(k_t \sin \theta_t x + k_t \cos \theta_t z)} \end{cases}$$

$$\begin{cases} \underline{\hat{H}}_i = \frac{\hat{E}_{i0}}{\eta_1} e^{-j(k_i \sin \theta_i x + k_i \cos \theta_i z)} \underline{a}_y \\ \underline{\hat{H}}_r = -\frac{\hat{E}_{r0}}{\eta_1} e^{-j(k_r \sin \theta_r x - k_r \cos \theta_r z)} \underline{a}_y \\ \underline{\hat{H}}_t = \frac{\hat{E}_{t0}}{\eta_2} e^{-j(k_t \sin \theta_t x + k_t \cos \theta_t z)} \underline{a}_y \end{cases}$$



Apply B.C., at $z=0$,

$$\underline{a}_x \cdot \underline{\hat{E}}_i \Big|_{z=0} + \underline{a}_x \cdot \underline{\hat{E}}_r \Big|_{z=0} = \underline{a}_x \cdot \underline{\hat{E}}_t \Big|_{z=0}$$

$$\underline{a}_y \cdot \underline{\hat{H}}_i \Big|_{z=0} + \underline{a}_y \cdot \underline{\hat{H}}_r \Big|_{z=0} = \underline{a}_y \cdot \underline{\hat{H}}_t \Big|_{z=0}$$

$$\begin{cases} (\hat{E}_{i0} + \hat{E}_{r0}) \cos \theta_i = \hat{E}_{t0} \cos \theta_t \\ \frac{\hat{E}_{i0}}{\eta_1} - \frac{\hat{E}_{r0}}{\eta_1} = \frac{\hat{E}_{t0}}{\eta_2} \end{cases} \qquad \begin{cases} \frac{\hat{E}_{t0}}{\hat{E}_{i0}} = \left(1 + \frac{\hat{E}_{r0}}{\hat{E}_{i0}} \right) \frac{\cos \theta_i}{\cos \theta_t} \\ \frac{1}{\eta_1} - \frac{1}{\eta_1} \left(\frac{\hat{E}_{r0}}{\hat{E}_{i0}} \right) = \frac{1}{\eta_2} \left(\frac{\hat{E}_{t0}}{\hat{E}_{i0}} \right) \end{cases}$$

$$\text{Reflection coefficient } \Gamma_{//} \equiv \frac{\hat{E}_{r0}}{\hat{E}_{i0}} = \frac{\eta_2 \cos \theta_t - \eta_1 \cos \theta_i}{\eta_2 \cos \theta_t + \eta_1 \cos \theta_i} \Rightarrow \Gamma_{//} = \frac{n_1 \cos \theta_t - n_2 \cos \theta_i}{n_1 \cos \theta_t + n_2 \cos \theta_i}$$

$$\begin{aligned}
 \Gamma_{//} &= \frac{-\sqrt{\mu_1/\epsilon_1} \cos \theta_i + \sqrt{\mu_2/\epsilon_2} \cos \theta_t}{\sqrt{\mu_1/\epsilon_1} \cos \theta_i + \sqrt{\mu_2/\epsilon_2} \cos \theta_t} = \frac{\cos \theta_t - \sqrt{\epsilon_{2r}/\epsilon_{1r}} \cos \theta_i}{\cos \theta_i + \sqrt{\epsilon_{2r}/\epsilon_{1r}} \cos \theta_i} \\
 &= \frac{\sqrt{1 - \sin^2 \theta_t} - \sqrt{\epsilon_{2r}/\epsilon_{1r}} \cos \theta_i}{\sqrt{1 - \sin^2 \theta_t} + \sqrt{\epsilon_{2r}/\epsilon_{1r}} \cos \theta_i} = \frac{\sqrt{1 - (\epsilon_{1r}/\epsilon_{2r}) \sin^2 \theta_i} - \sqrt{\epsilon_{2r}/\epsilon_{1r}} \cos \theta_i}{\sqrt{1 - (\epsilon_{1r}/\epsilon_{2r}) \sin^2 \theta_i} + \sqrt{\epsilon_{2r}/\epsilon_{1r}} \cos \theta_i} \\
 &= \frac{-\cos \theta_i + \sqrt{\epsilon_{1r}/\epsilon_{2r}} \sqrt{1 - (\epsilon_{1r}/\epsilon_{2r}) \sin^2 \theta_i}}{\cos \theta_i + \sqrt{\epsilon_{1r}/\epsilon_{2r}} \sqrt{1 - (\epsilon_{1r}/\epsilon_{2r}) \sin^2 \theta_i}} = \frac{-\cos \theta_i + \sqrt{(\epsilon_{1r}/\epsilon_{2r}) - (\epsilon_{1r}/\epsilon_{2r})^2 \sin^2 \theta_i}}{\cos \theta_i + \sqrt{(\epsilon_{1r}/\epsilon_{2r}) - (\epsilon_{1r}/\epsilon_{2r})^2 \sin^2 \theta_i}} \\
 &= \frac{-(\epsilon_{2r}/\epsilon_{1r}) \cos \theta_i + \sqrt{(\epsilon_{2r}/\epsilon_{1r}) - \sin^2 \theta_i}}{(\epsilon_{2r}/\epsilon_{1r}) \cos \theta_i + \sqrt{(\epsilon_{2r}/\epsilon_{1r}) - \sin^2 \theta_i}} = \frac{-(n_2/n_1)^2 \cos \theta_i + \sqrt{(n_2^2/n_1^2) - \sin^2 \theta_i}}{(n_2/n_1)^2 \cos \theta_i + \sqrt{(n_2^2/n_1^2) - \sin^2 \theta_i}}
 \end{aligned}$$

$$\Gamma_{//} = \frac{-\cos \theta_i + \sqrt{(\epsilon_{1r}/\epsilon_{2r})} \sqrt{(\epsilon_{2r}/\epsilon_{1r}) - \sin^2 \theta_i}}{\cos \theta_i + \sqrt{(\epsilon_{1r}/\epsilon_{2r})} \sqrt{(\epsilon_{2r}/\epsilon_{1r}) - \sin^2 \theta_i}} \Rightarrow \Gamma_{//} = \frac{-(n_2/n_1)^2 \cos \theta_i + \sqrt{(n_2^2/n_1^2) - \sin^2 \theta_i}}{(n_2/n_1)^2 \cos \theta_i + \sqrt{(n_2^2/n_1^2) - \sin^2 \theta_i}}$$

Transmission coefficient $\mathfrak{T}_{//} = \frac{\hat{E}_{t0}}{\hat{E}_{i0}} = \frac{2\eta_2 \cos \theta_i}{\eta_1 \cos \theta_i + \eta_2 \cos \theta_t} \Rightarrow \mathfrak{T}_{//} = \frac{2n_1 \cos \theta_i}{n_2 \cos \theta_i + n_1 \cos \theta_t}$

$$\begin{aligned}
 \mathfrak{T}_{//} &= \frac{2 \cos \theta_i}{\cos \theta_i + \sqrt{\epsilon_{2r}/\epsilon_{1r}} \cos \theta_t} = \frac{2 \cos \theta_i}{\sqrt{1 - \sin^2 \theta_t} + \sqrt{\epsilon_{2r}/\epsilon_{1r}} \cos \theta_t} \\
 &= \frac{2 \cos \theta_i}{\sqrt{\epsilon_{2r}/\epsilon_{1r}} \cos \theta_i + \sqrt{1 - (\epsilon_{1r}/\epsilon_{2r}) \sin^2 \theta_i}} = \frac{2\sqrt{(\epsilon_{1r}/\epsilon_{2r})} \cos \theta_i}{\cos \theta_i + \sqrt{(\epsilon_{1r}/\epsilon_{2r})} \sqrt{1 - (\epsilon_{1r}/\epsilon_{2r}) \sin^2 \theta_i}} \\
 &= \frac{2\sqrt{(\epsilon_{2r}/\epsilon_{1r})} \cos \theta_i}{(\epsilon_{2r}/\epsilon_{1r}) \cos \theta_i + \sqrt{(\epsilon_{2r}/\epsilon_{1r}) - \sin^2 \theta_i}} = \frac{2(n_2/n_1) \cos \theta_i}{(n_2^2/n_1^2) \cos \theta_i + \sqrt{(n_2^2/n_1^2) - \sin^2 \theta_i}}
 \end{aligned}$$

$$\mathfrak{T}_{//} = \frac{2\sqrt{(\epsilon_{1r}/\epsilon_{2r})} \cos \theta_i}{\cos \theta_i + \sqrt{(\epsilon_{1r}/\epsilon_{2r})} \sqrt{1 - (\epsilon_{1r}/\epsilon_{2r}) \sin^2 \theta_i}} \Rightarrow \mathfrak{T}_{//} = \frac{2(n_2/n_1) \cos \theta_i}{(n_2^2/n_1^2) \cos \theta_i + \sqrt{(n_2^2/n_1^2) - \sin^2 \theta_i}}$$

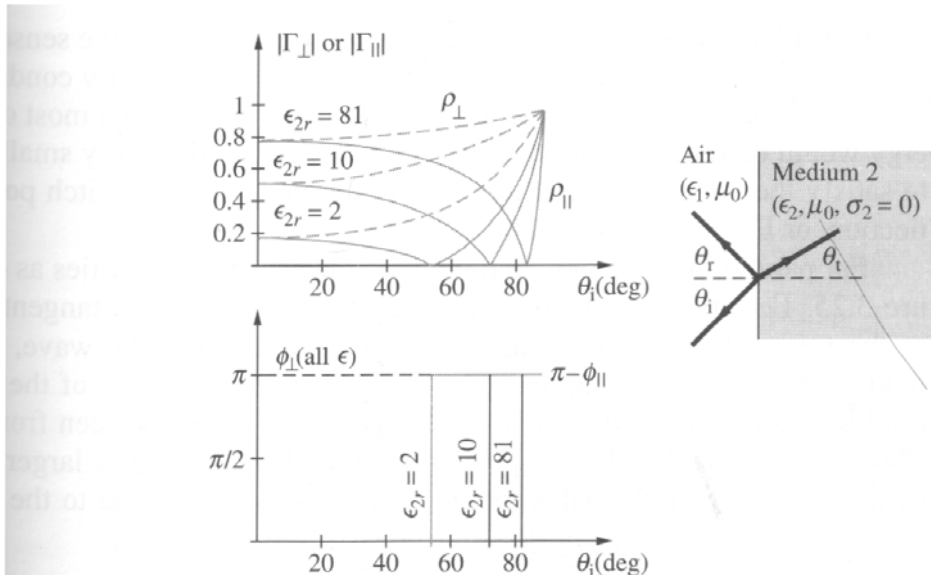


FIGURE 3.30. Reflection coefficient versus the angle of incidence. (a) Magnitude and phase of reflection coefficient for perpendicular (Γ_{\perp}) and parallel (Γ_{\parallel}) polarization versus angle of incidence θ_i for distilled water ($\epsilon_{2r} = 81$), flint glass ($\epsilon_{2r} = 10$) and paraffin ($\epsilon_{2r} = 2$), all assumed to be lossless. For clarity, the complement of the phase angle ϕ_{\parallel} (i.e., $\pi - \phi_{\parallel}$) is sketched, rather than ϕ_{\parallel} itself. In all cases, $\phi_{\parallel} = \pi$ for $\theta_i < \theta_{iB}$, and $\phi_{\parallel} = 0$ for $\theta_i > \theta_{iB}$, while $\phi_{\perp} = \pi$ for all θ_i .

$$1 + \Gamma_{\parallel} = \mathfrak{I}_{\parallel} \left(\frac{\cos \theta_t}{\cos \theta_i} \right)$$

Brewster's angle θ_{iB}

If $r_{\parallel} = 0$ (total transmission), then $n_1 \cos \theta_t - n_2 \cos \theta_i = 0$

$$\Rightarrow \cos \theta_t = (n_2/n_1) \cos \theta_i$$

By Snell's law of refraction, $n_1 \sin \theta_i = n_2 \sin \theta_t$

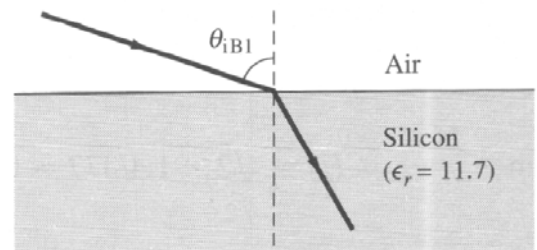
$$\Rightarrow \sin \theta_t = (n_1/n_2) \sin \theta_i$$

$$\begin{aligned} \therefore \sin^2 \theta_t + \cos^2 \theta_t &= \left(\frac{n_1}{n_2} \right)^2 \sin^2 \theta_i + \left(\frac{n_2}{n_1} \right)^2 \cos^2 \theta_i \\ &= 1 = \sin^2 \theta_i + \cos^2 \theta_i \end{aligned}$$

$$\left[1 - \left(\frac{n_1}{n_2} \right)^2 \right] \sin^2 \theta_i = \left[\left(\frac{n_2}{n_1} \right)^2 - 1 \right] \cos^2 \theta_i \quad \text{(or)} \quad \left(1 - \frac{\epsilon_1}{\epsilon_2} \right) \sin^2 \theta_i = \left(\frac{\epsilon_2}{\epsilon_1} - 1 \right) \cos^2 \theta_i$$

$$\tan^2 \theta_i = \left(\frac{n_2}{n_1} \right)^2 \Rightarrow \theta_{iB} = \tan^{-1} \frac{n_2}{n_1}$$

for an air-silicon ($\epsilon_r = 11.7$) interface



$$\theta_{iB} = \tan^{-1} \sqrt{\epsilon_r} = \tan^{-1} \sqrt{11.7} \approx 73.7^\circ$$

Normal incidence on a dielectric boundary

$$\theta_i = \theta_r = \theta_t = 0^\circ$$

$$\begin{cases} \underline{\hat{E}}_i = \underline{\hat{E}}_{i0} e^{-jk_i z} \underline{a}_x \\ \underline{\hat{E}}_r = \underline{\hat{E}}_{r0} e^{-jk_i z} \underline{a}_x \\ \underline{\hat{E}}_t = \underline{\hat{E}}_{t0} e^{-jk_t z} \underline{a}_x \end{cases} \quad \begin{cases} \underline{\hat{H}}_i = \frac{\underline{\hat{E}}_{i0}}{\eta_1} e^{-jk_i z} \underline{a}_y \\ \underline{\hat{H}}_r = -\frac{\underline{\hat{E}}_{r0}}{\eta_1} e^{-jk_i z} \underline{a}_y \\ \underline{\hat{H}}_t = \frac{\underline{\hat{E}}_{t0}}{\eta_2} e^{-jk_t z} \underline{a}_y \end{cases}$$

$$\text{Reflection coefficient } \Gamma_{11} \equiv \frac{\underline{\hat{E}}_{r0}}{\underline{\hat{E}}_{i0}} = \frac{\eta_2 - \eta_1}{\eta_2 + \eta_1} = \Gamma_{\perp}$$

$$\text{Transmission coefficient } \mathfrak{T}_{11} \equiv \frac{\underline{\hat{E}}_{t0}}{\underline{\hat{E}}_{i0}} = \frac{2\eta_2}{\eta_2 + \eta_1} = \mathfrak{T}_{\perp}$$

In most cases, $\mu_1 = \mu_2 = \mu_0$

$$\Gamma = \frac{\sqrt{\epsilon_1} - \sqrt{\epsilon_2}}{\sqrt{\epsilon_1} + \sqrt{\epsilon_2}} = \frac{n_1 - n_2}{n_1 + n_2} \quad \mathfrak{T} = \frac{2\sqrt{\epsilon_1}}{\sqrt{\epsilon_1} + \sqrt{\epsilon_2}} = \frac{2n_1}{n_1 + n_2}$$

$$\begin{aligned} \underline{\hat{E}}_1 &= \underline{\hat{E}}_i + \underline{\hat{E}}_r = \underline{\hat{E}}_{i0} e^{-jk_i z} \underline{a}_x + \underline{\hat{E}}_{r0} e^{jk_i z} \underline{a}_x = \underline{\hat{E}}_{i0} e^{-jk_i z} \underline{a}_x + \Gamma \underline{\hat{E}}_{i0} e^{jk_i z} \underline{a}_x \\ &= \underline{\hat{E}}_{i0} e^{-jk_i z} (1 + \Gamma e^{j2k_i z}) \underline{a}_x \end{aligned}$$

$$\begin{aligned} \underline{\hat{H}}_1 &= \underline{\hat{H}}_i + \underline{\hat{H}}_r = \frac{\underline{\hat{E}}_{i0}}{\eta_1} e^{-jk_i z} \underline{a}_y - \frac{\underline{\hat{E}}_{r0}}{\eta_1} e^{jk_i z} \underline{a}_y = \frac{\underline{\hat{E}}_{i0}}{\eta_1} e^{-jk_i z} \underline{a}_y - \Gamma \frac{\underline{\hat{E}}_{i0}}{\eta_1} e^{jk_i z} \underline{a}_y \\ &= \frac{\underline{\hat{E}}_{i0}}{\eta_1} e^{-jk_i z} (1 - \Gamma e^{j2k_i z}) \underline{a}_y \end{aligned}$$

$$|\underline{\hat{E}}_1| = |\underline{\hat{E}}_{i0}| |1 + \Gamma e^{j2k_i z}|$$

for $\Gamma > 0$

$$|\underline{\hat{E}}_1|_{\max} = |\underline{\hat{E}}_{i0}| (1 + \Gamma)$$

maximum occurs when

$$2k_1 z = -2n\pi \quad n = 0, 1, 2, \dots$$

$$k_1 = 2\pi / \lambda_1 \quad z = -n\lambda_1 / 2$$

$$|\underline{\hat{E}}_1|_{\min} = |\underline{\hat{E}}_{i0}| (1 - \Gamma)$$

minima occurs when $2k_1 z = -m\pi \quad m = 1, 3, 5, \dots$

$$z = -m\lambda_1 / 4$$

for $\Gamma < 0$

$$\left| \hat{E}_1 \right|_{\max} = \left| \hat{E}_{i0} \right| (1 - \Gamma)$$

$$\left| \hat{E}_1 \right|_{\min} = \left| \hat{E}_{i0} \right| (1 + \Gamma)$$

$$\therefore \begin{cases} \left| \hat{E}_1 \right|_{\max} = \left| \hat{E}_{i0} \right| (1 + |\Gamma|) \\ \left| \hat{E}_1 \right|_{\min} = \left| \hat{E}_{i0} \right| (1 - |\Gamma|) \end{cases}$$

$$S = \frac{\left| \hat{E}_1 \right|_{\max}}{\left| \hat{E}_1 \right|_{\min}} = \frac{1 + |\Gamma|}{1 - |\Gamma|} \quad |\Gamma| = \frac{S - 1}{S + 1}$$

Oblique incidence on a perfect conductor

Incident wave	Reflected wave
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$$\underline{\hat{E}}_i = \hat{E}_{i0} e^{-jk_i \cdot \underline{r}} (\underline{a}_y \times \underline{a}_{ni})$$

$$\underline{\hat{E}}_r = \hat{E}_{r0} e^{-jk_r \cdot \underline{r}} (\underline{a}_y \times \underline{a}_{nr})$$

$$\underline{\hat{H}}_i = \frac{\hat{E}_{i0} e^{-jk_i \cdot \underline{r}}}{\eta_1} \underline{a}_y$$

$$\underline{\hat{H}}_r = \frac{\hat{E}_{r0} e^{-jk_r \cdot \underline{r}}}{\eta_1} \underline{a}_y$$

$$\underline{a}_{ni} = \sin \theta_i \underline{a}_x + \cos \theta_i \underline{a}_z$$

$$\underline{a}_{nr} = \sin \theta_r \underline{a}_x - \cos \theta_r \underline{a}_z$$

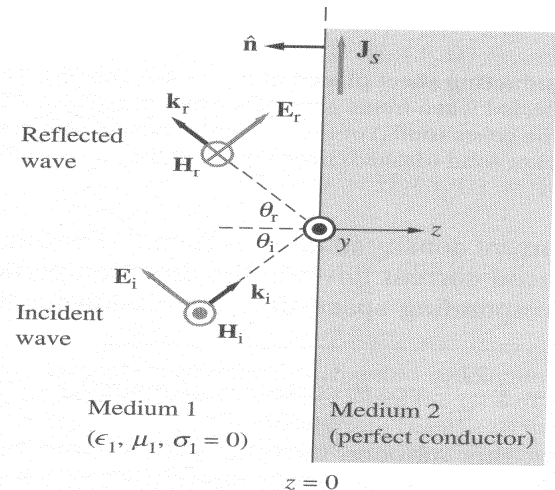
$$\underline{k}_i = k_i \sin \theta_i \underline{a}_x + k_i \cos \theta_i \underline{a}_z$$

$$\underline{k}_r = k_r \sin \theta_r \underline{a}_x - k_r \cos \theta_r \underline{a}_z$$

$$\begin{cases} \underline{\hat{E}}_i = \hat{E}_{i0} (\cos \theta_i \underline{a}_x - \sin \theta_i \underline{a}_z) e^{-j(k_i \sin \theta_i x + k_i \cos \theta_i z)} \\ \underline{\hat{E}}_r = \hat{E}_{r0} (\cos \theta_r \underline{a}_x + \sin \theta_r \underline{a}_z) e^{-j(k_r \sin \theta_r x - k_r \cos \theta_r z)} \end{cases}$$

$$\begin{cases} \underline{\hat{H}}_i = \frac{\hat{E}_{i0}}{\eta_1} e^{-j(k_i \sin \theta_i x + k_i \cos \theta_i z)} \underline{a}_y \\ \underline{\hat{H}}_r = -\frac{\hat{E}_{r0}}{\eta_1} e^{-j(k_r \sin \theta_r x - k_r \cos \theta_r z)} \underline{a}_y \end{cases}$$

Apply B.C., at $z=0$,



$$\underline{a}_x \cdot \underline{\hat{E}}_i|_{z=0} + \underline{a}_x \cdot \underline{\hat{E}}_r|_{z=0} = 0$$

$$\underline{a}_y \cdot \underline{\hat{H}}_i|_{z=0} + \underline{a}_y \cdot \underline{\hat{H}}_r|_{z=0} = J_s$$

$$\hat{E}_{i0} = -\hat{E}_{r0}, \quad \theta_i = \theta_r$$

$$\begin{aligned} \underline{\hat{E}}_1 &= \underline{\hat{E}}_i + \underline{\hat{E}}_r = \hat{E}_{i0} \cos \theta_i \left(e^{-jk_1 \cos \theta_i z} - e^{+jk_1 \cos \theta_i z} \right) e^{-jk_1 \sin \theta_i x} \underline{a}_x - \hat{E}_{i0} \sin \theta_i \left(e^{-jk_1 \cos \theta_i z} + e^{+jk_1 \cos \theta_i z} \right) e^{-jk_1 \sin \theta_i x} \underline{a}_z \\ &= -2\hat{E}_{i0} \left[\cos \theta_i \sin(k_1 z \cos \theta_i) \underline{a}_x - \sin \theta_i \cos(k_1 z \cos \theta_i) \underline{a}_z \right] e^{-jk_1 \sin \theta_i x} \end{aligned}$$

$$\underline{\hat{H}}_1 = \underline{\hat{H}}_i + \underline{\hat{H}}_r = \frac{2\hat{E}_{i0}}{\eta_1} \cos(k_1 z \cos \theta_i) e^{-jk_1 \sin \theta_i x} \underline{a}_y$$

Time-average Poynting vector

$$\begin{aligned} \underline{S}_{av} &= \frac{1}{2} \text{Re} \left\{ \underline{\hat{E}}_1 \times \underline{\hat{H}}_1^* \right\} \\ &= \frac{2\hat{E}_{i0}^2}{\eta_1} \sin \theta_i \cos^2(k_1 \cos \theta_i z) \underline{a}_x \end{aligned}$$

at $z = 0$, induced current in the conductor

$$\underline{J}_s = \underline{a}_n \times \underline{H}_1|_{z=0} = \frac{2\hat{E}_{i0}}{\eta_1} e^{-jk_1 \sin \theta_i x} \underline{a}_y$$

Normal incidence on a perfect conductor

$$\theta_i = \theta_r = 0$$

$$\underline{\hat{E}}_1 = j2\hat{E}_{i0} \sin(k_1 z) \underline{a}_x$$

$$\underline{\hat{H}}_1 = -\frac{2\hat{E}_{i0}}{\eta_1} \cos(k_1 z) \underline{a}_y$$

$$\begin{aligned} \underline{S}_{av} &= \frac{1}{2} \text{Re} \left\{ \underline{\hat{E}}_1 \times \underline{\hat{H}}_1^* \right\} = \frac{1}{2} \text{Re} \left\{ \left[j2\hat{E}_{i0} \sin(k_1 z) \right] \underline{a}_x \times \left[-\frac{2\hat{E}_{i0}^*}{\eta_1} \cos(k_1 z) \right] \underline{a}_y \right\} \\ &= \frac{1}{2} \text{Re} \left\{ \left[-j \frac{2|\hat{E}_{i0}|^2}{\eta_1} \sin(2k_1 z) \right] \underline{a}_z \right\} = 0 \end{aligned}$$

at $z = 0$, induced current in the conductor

$$\underline{J}_s = \underline{a}_n \times \underline{H}_1|_{z=0} = \left[(-\underline{a}_z) \times (\underline{a}_y) \right] \frac{2\hat{E}_{i0}}{\eta_1} = \frac{2\hat{E}_{i0}}{\eta_1} \underline{a}_x$$

Total internal reflection

No refraction $\Rightarrow \theta_t = 90^\circ$

$\Rightarrow n_1 \sin \theta_{ic} = n_2$

$\Rightarrow \sin \theta_{ic} = \sqrt{\frac{\epsilon_2}{\epsilon_1}} = \frac{n_2}{n_1}$

\Rightarrow critical angle $\theta_{ic} = \sin^{-1} \sqrt{\frac{\epsilon_2}{\epsilon_1}} = \sin^{-1} \frac{n_2}{n_1}$

if $\theta_i > \theta_{ic} \Rightarrow$ total internal reflection

$\Rightarrow \frac{n_1}{n_2} \sin \theta_i > 1 \quad \left(\because \frac{n_1}{n_2} \sin \theta_{ic} = 1 \right)$

$\Rightarrow \frac{n_1}{n_2} \sin \theta_i = \sin \theta_t > 1 \quad \Rightarrow \quad \theta_t$ imaginary (if θ_t is real $\Rightarrow \sin \theta_t \leq 1$)

$\Rightarrow \cos \theta_t = \sqrt{1 - \sin^2 \theta_t} = \pm j \sqrt{\frac{n_1^2}{n_2^2} \sin^2 \theta_i - 1}$ purely imaginary

$\Rightarrow \cos \theta_t = -j \sqrt{\frac{n_1^2}{n_2^2} \sin^2 \theta_i - 1} = -j \sqrt{\frac{\epsilon_{1r}}{\epsilon_{2r}} \sin^2 \theta_i - 1}$ ensuring the refracted EM field in medium 2

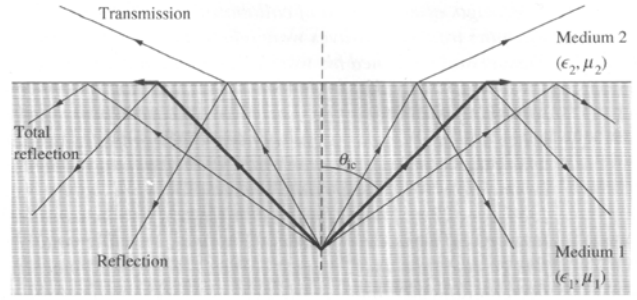
attenuates with increasing distance from the boundary.

(IEEE convention, $e^{j(\omega t - kz)} \Rightarrow e^{-jkz} \Rightarrow (-j) \cdot (-j) = -1 \Rightarrow$

$$\begin{aligned} \Gamma_{\perp} &= \frac{n_i \cos \theta_i - n_t \cos \theta_t}{n_i \cos \theta_i + n_t \cos \theta_t} = \frac{\cos \theta_i + j \sqrt{\frac{\epsilon_{2r}}{\epsilon_{1r}} \sqrt{\frac{\epsilon_{1r}}{\epsilon_{2r}} \sin^2 \theta_i - 1}}}{\cos \theta_i - j \sqrt{\frac{\epsilon_{2r}}{\epsilon_{1r}} \sqrt{\frac{\epsilon_{1r}}{\epsilon_{2r}} \sin^2 \theta_i - 1}}} = \frac{\cos \theta_i + j \sqrt{\sin^2 \theta_i - \frac{\epsilon_{2r}}{\epsilon_{1r}}}}{\cos \theta_i - j \sqrt{\sin^2 \theta_i - \frac{\epsilon_{2r}}{\epsilon_{1r}}}} \\ &= \frac{\sqrt{\epsilon_{1r}} \cos \theta_i + j \sqrt{\epsilon_{1r} \sin^2 \theta_i - \epsilon_{2r}}}{\sqrt{\epsilon_{1r}} \cos \theta_i - j \sqrt{\epsilon_{1r} \sin^2 \theta_i - \epsilon_{2r}}} = \frac{n_1 \cos \theta_i + j \sqrt{n_1^2 \sin^2 \theta_i - n_2^2}}{n_1 \cos \theta_i - j \sqrt{n_1^2 \sin^2 \theta_i - n_2^2}} \end{aligned}$$

(Optics convention, $e^{-j(\omega t - kz)} \Rightarrow e^{jkz} \Rightarrow j:j = -1$)

$$\Gamma_{\perp} = \frac{\cos \theta_i - i \sqrt{\sin^2 \theta_i - \frac{\epsilon_{2r}}{\epsilon_{1r}}}}{\cos \theta_i + i \sqrt{\sin^2 \theta_i - \frac{\epsilon_{2r}}{\epsilon_{1r}}}} = \frac{\cos \theta_i - i \sqrt{\sin^2 \theta_i - \frac{n_2^2}{n_1^2}}}{\cos \theta_i + i \sqrt{\sin^2 \theta_i - \frac{n_2^2}{n_1^2}}} = \frac{n_1^2 \cos \theta_i - i \sqrt{n_1^2 \sin^2 \theta_i - n_2^2}}{n_1^2 \cos \theta_i + i \sqrt{n_1^2 \sin^2 \theta_i - n_2^2}}$$



$$\Gamma_{\perp} = \frac{\cos \theta_i + j\sqrt{\sin^2 \theta_i - \varepsilon_{21}}}{\cos \theta_i - j\sqrt{\sin^2 \theta_i - \varepsilon_{21}}} = e^{j\phi_{\perp}} \quad \varepsilon_{21} = \frac{\varepsilon_{2r}}{\varepsilon_{1r}}$$

$$e^{j\phi} = z / z^* = e^{j2\delta}, \quad z = e^{j\delta}$$

$$\Rightarrow \tan \frac{\phi_{\perp}}{2} = \frac{\sqrt{\sin^2 \theta_i - \varepsilon_{21}}}{\cos \theta_i}$$

(IEEE convention)

$$\begin{aligned} \Gamma_{//} &= \frac{n_t \cos \theta_i - n_r \cos \theta_r}{n_i \cos \theta_i + n_r \cos \theta_r} = -\frac{\sqrt{\frac{\varepsilon_{2r}}{\varepsilon_{1r}} \cos \theta_i + j\sqrt{\frac{\varepsilon_{1r}}{\varepsilon_{2r}} \sin^2 \theta_i - 1}}{\sqrt{\frac{\varepsilon_{2r}}{\varepsilon_{1r}} \cos \theta_i - j\sqrt{\frac{\varepsilon_{1r}}{\varepsilon_{2r}} \sin^2 \theta_i - 1}}} = -\frac{\frac{\varepsilon_{2r}}{\varepsilon_{1r}} \cos \theta_i + j\sqrt{\sin^2 \theta_i - \frac{\varepsilon_{2r}}{\varepsilon_{1r}}}}{\frac{\varepsilon_{2r}}{\varepsilon_{1r}} \cos \theta_i - j\sqrt{\sin^2 \theta_i - \frac{\varepsilon_{2r}}{\varepsilon_{1r}}}} = e^{j\phi_{//}} \\ &= \frac{\varepsilon_{2r} \cos \theta_i + j\sqrt{\varepsilon_{1r}} \sqrt{\varepsilon_{1r} \sin^2 \theta_i - \varepsilon_{2r}}}{\varepsilon_{2r} \cos \theta_i - j\sqrt{\varepsilon_{1r}} \sqrt{\varepsilon_{1r} \sin^2 \theta_i - \varepsilon_{2r}}} = \frac{n_2^2 \cos \theta_i + jn_1 \sqrt{n_1^2 \sin^2 \theta_i - n_2^2}}{n_2^2 \cos \theta_i - jn_1 \sqrt{n_1^2 \sin^2 \theta_i - n_2^2}} \end{aligned}$$

(Optics convention)

$$\Gamma_{//} = \frac{n_2^2 \cos \theta_i - in_1 \sqrt{n_1^2 \sin^2 \theta_i - n_2^2}}{n_2^2 \cos \theta_i + in_1 \sqrt{n_1^2 \sin^2 \theta_i - n_2^2}}$$

$$\Gamma_{//} = -\frac{\varepsilon_{21} \cos \theta_i + j\sqrt{\sin^2 \theta_i - \varepsilon_{21}}}{\varepsilon_{21} \cos \theta_i - j\sqrt{\sin^2 \theta_i - \varepsilon_{21}}} = e^{j\phi_{//}} \quad \Rightarrow \quad \tan \frac{\phi_{//}}{2} = \frac{\sqrt{\sin^2 \theta_i - \varepsilon_{21}}}{\varepsilon_{21} \cos \theta_i}$$

linearly polarized wave \rightarrow circularly or elliptically polarized wave by using phase change in total reflection \Rightarrow linearly polarized wave incident at an angle of 45° with respect to the plane of

incidence $\Rightarrow |E_{//}| = |E_{\perp}|$

$$\Delta\phi = \phi_{\perp} - \phi_{//} \quad \Rightarrow \quad \tan \frac{\Delta\phi}{2} = \frac{\tan(\phi_{\perp}/2) - \tan(\phi_{//}/2)}{1 + \tan(\phi_{\perp}/2)\tan(\phi_{//}/2)} = \frac{\cos \theta_i \sqrt{\sin^2 \theta_i - \varepsilon_{21}}}{\sin^2 \theta_i}$$

$$\hat{Z}_s = \frac{\hat{E}_t}{\hat{J}_s} = \frac{\hat{E}_y}{\hat{H}_z} =$$